

September 11, 2012

Dominant $1/m_c$ Contribution To The Mass Difference in D^0 - \overline{D}^0 Mixing

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We re-visit the problem of non-perturbative contribution to the mass difference in $D^0 - \overline{D}^0$ mixing within the Standard Model (SM). There are two known approaches to take this contribution into account. In the long-distance approach one analyzes contribution to the $D^0 - \overline{D}^0$ from exclusive channels due to production of intermediate charmless hadronic states. In the short-distance approach one employs the Operator Product Expansion (OPE) techniques with the non-perturbative contribution coming from the diagrams containing low energy-momentum quark states - quark-antiquark condensates. Within the latter approach, the non-perturbative contribution to the normalized mass difference in $D^0 - \overline{D}^0$ mixing, $x_D = \Delta M_D / \Gamma_D$, is estimated to be $\text{few} \times 10^{-3}$, that is to say very close to the experimental value of x_D . We attempt to verify quantitatively this estimate here. We find that the actual prediction of the short-distance approach is few times less than the estimate above and order of magnitude less than the experimental value of x_D .

PRESENTED AT

The Fifth International Workshop On Charm Physics
Honolulu, Hawaii, May 13–17, 2012

$D^0 - \overline{D}^0$ oscillations may serve as an efficient tool for indirect searches for New Physics (NP). Recent discovery of CP-violation in the charm quark involving processes [1, 2] with the experimental data possibly exceeding the SM predictions has revived the interest to the charm sector as the place where a signal for New Physics (NP) may be detected.

Within the perturbative approach, the Standard Model contribution to the mass and the width difference in $D^0 - \overline{D}^0$ mixing is suppressed by several orders of magnitude [3, 4] as compared to the experimental values of these quantities. The experimental values of the normalized mass and width differences in $D^0 - \overline{D}^0$ mixing [5],

$$x_D^{exp} = \frac{\Delta M_D^{exp}}{\Gamma_D} = (6.3 \pm 1.9) \times 10^{-3} \quad (1)$$

$$y_D^{exp} = \frac{\Delta \Gamma_D^{exp}}{2\Gamma_D} = (7.5 \pm 1.2) \times 10^{-3} \quad (2)$$

may be either due to non-perturbative effects within the Standard Model [6, 7, 8, 9, 10] or due to New Physics effects (see [11, 12, 13] and references therein). Whether one of these two effects dominates or (1) and/or (2) are due to interference of these two effects, is still unclear mainly because of the lack of precise evaluation of the non-perturbative SM contribution.

In this paper the problem of non-perturbative contribution to the mass difference in $D^0 - \overline{D}^0$ mixing within the Standard Model is readdressed.

Two ways of taking into account the non-perturbative contribution to $D^0 - \overline{D}^0$ mixing are presently known. The long-distance approach considers the exclusive channels of $D^0 - \overline{D}^0$ mixing, when this process occurs due to production of intermediate charmless hadronic states. It has been estimated using this approach that $x_D^{SM} \sim 10^{-2}$ and $y_D^{SM} \sim 10^{-2}$.

Another (short-distance) approach is based on the use of Operator Product Expansion (OPE) techniques with the non-perturbative effects coming from the diagrams containing low energy-momentum quark states - quark-antiquark condensates. The contribution of these diagrams corresponds to higher-order (as compared to the dimension-six four-fermion operators) $1/m_c$ terms in the OPE. The estimates done using this approach predict [8] $x_D^{SM} \sim \text{few} \times 10^{-3}$, and $y_D \sim 10^{-3}$. Apparent discrepancy between the results of the two approaches may be both because of the Operator Product Expansion failure (as μ_{had}/m_c is not a small quantity), and because of possible cancelations in the contribution of different exclusive channels.

In this paper we attempt to evaluate quantitatively the non-perturbative contribution to the mass difference in $D^0 - \overline{D}^0$ mixing using the OPE techniques*. Our preliminary results show that the actual prediction of the short-distance approach

*Quantitative evaluation of subleading diagrams for the width difference has been done in [14].

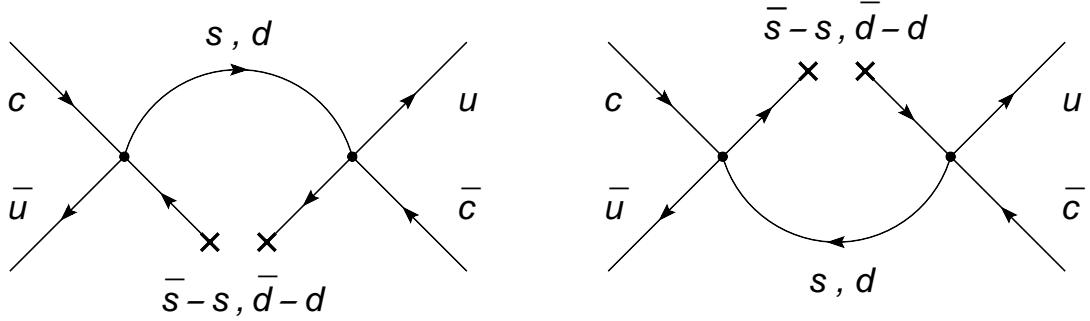


Figure 1: Diagrams with two hanging quark lines.

is few times less than the estimate of ref. [8] and order of magnitude less than the experimental value of x_D .

Within the Standard Model, $D^0 - \bar{D}^0$ oscillations are analyzed using the two generation mixing approach (neglecting the contribution of the diagrams with b-quark intermediate states - this approach is well justified (at least when the mass difference is of the interest) due to the suppressed mixing of the third generation of quarks with the first two). With this assumption, $D^0 - \bar{D}^0$ mixing vanishes in the well-defined theoretical limit of exact flavor $SU(3)$ symmetry. In the real world, where the flavor $SU(3)$ is broken, the SM contribution to the $D^0 - \bar{D}^0$ is of course non-vanishing. Yet, it is suppressed by a certain power of the strange-to-charm mass ratio. In particular, to the lowest order in the perturbation theory, $x_D^{LO} \propto m_s^4/m_c^4$ and $y_D^{LO} \propto m_s^6/m_c^6$. More precisely, for the mass difference [4],

$$(x_D)^{LO} = \frac{G_F^2 m_c^2 f_D^2 M_D}{3\pi^2 \Gamma_D} \left(\frac{m_s^4}{m_c^4} \right) \sin^2 \theta_C \cos^2 \theta_C \times \left\{ \frac{5}{4} [C_2^2 - 2C_1C_2 - 3C_1^2] \bar{B}_D^{(S)} - C_2^2 B_D \right\} \approx 1.7 \times 10^{-6} \quad (3)$$

Thus, within the SM, to the lowest order in the perturbation theory x_D is smaller by several orders of magnitude than its experimental value.

It has been argued however [6, 7, 8] that the SM contribution to x_D may be significantly greater due to the diagrams involving interactions of the valence charm and up quarks and antiquarks with (non-perturbative) low energy-momentum intermediate QCD background quark states (light quark-antiquark condensates). The simplest case with a single low energy-momentum quark state and the energy being transferred from D^0 to \bar{D}^0 (and back) by a single light quark propagator is depicted in Figure 1. The non-perturbative quark-antiquark state is denoted by a pair of hanging light quark lines.

It is also possible that all the intermediate light quark states are low energy-momentum states and the energy from D^0 to \bar{D}^0 and back is transferred by a hard

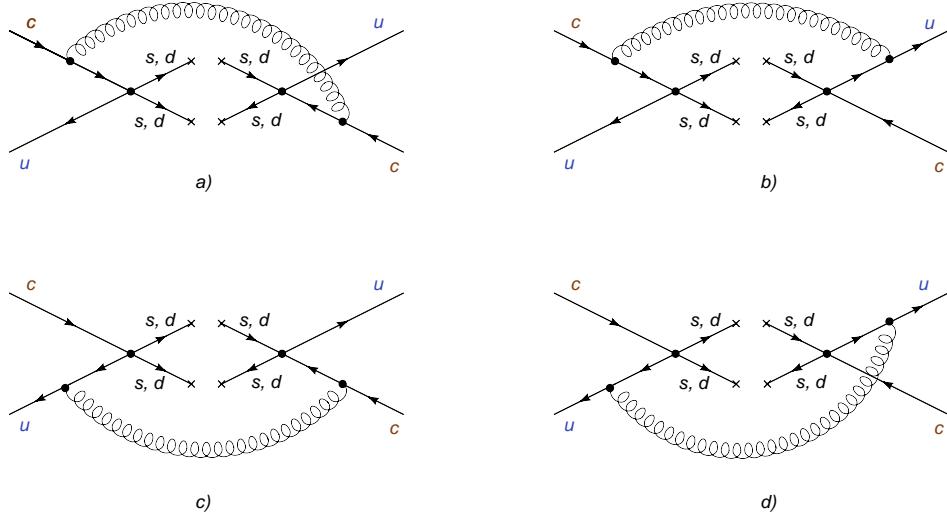


Figure 2: Examples of diagrams with four hanging light quark lines.

gluon. Examples of the diagrams with four hanging light quark lines (that describe such a scenario) are presented in Figure 2.

Note that each time one cuts a quark line, the strange or down quark propagator in the matrix element of $D^0 - \overline{D}^0$ transition must be replaced by a dimension-three product of strange or down quark operators. Then an extra factor of m_c^3 must appear in the denominator to preserve the dimensionality of the matrix element (and that of the mass and the width difference). Thus, the diagrams in Fig. 1 yield the matrix elements of $d = 9$ local operator products or $1/m_c^3$ suppressed terms in Operator Product Expansion. Subsequently, the diagrams in Fig. 2 yield the matrix elements of $d = 12$ local operator products or $1/m_c^6$ suppressed terms in Operator Product Expansion.

As the ratio of μ_{had}/m_c is not a small quantity, both the $1/m_c^3$ and the $1/m_c^6$ terms in OPE may be non-negligible or even dominant provided that their suppression in powers of m_s/m_c is weaker than that of the lowest-order term. It has been proven [6] using a group-theoretical analysis that mass insertion in each light quark propagator produces a factor of m_s^2 , whereas the interaction of the valence charm and up quarks and antiquarks with a QCD background quark-antiquark state produces a factor of m_s only. Thus, one may expect a softer flavor $SU(3)$ suppression in the diagrams with hanging quarks!

The rule of the suppression power counting in the diagrams with hanging quarks states as follows [8]:

- Cutting a quark line, we pay the price of a power suppression $\sim \mu_{had}^3/m_c^3$. Yet, flavor $SU(3)$ suppression (in this fermion line) is m_s/μ_{had} only and there is no loop factor. Altogether we have the enhancement $\sim 4\pi^2\mu_{had}^2/(m_s m_c)$, which

may yield $(x_D)^{d=9} \sim 10^{-4}$.

- Subsequently, cutting two quark lines (and inserting a gluon propagator that transfers the energy from D^0 to \bar{D}^0 and back), produces an enhancement factor of $\sim 4\pi^2(4\pi\alpha_s)\mu_{had}^4/(m_s^2m_c^2)$ (with $\alpha_s = \alpha_s(m_c)$). This may yield $(x_D)^{d=12} \sim \text{few} \times 10^{-3}$.

The latter estimate however turns to be somewhat altered by some additional suppression factors that emerge in the quantitative evaluation of the diagrams with hanging quarks. We illustrate this on the subset of the diagrams in Fig. 2 where the gluon couples to the external (charm and up) quark lines. The more comprehensive evaluation of all relevant diagram with hanging quarks will be presented in our forthcoming publication [15].

For the diagrams in Fig. 2 one gets

$$(x_D)^{d=12}_{external} = \left(\frac{2G_F^2m_c^2}{9}\right)(4\pi\alpha_s)\frac{f_D^2M_D}{\Gamma_D}\left(\frac{m_s^2}{m_c^2}\right)\left(\frac{\Lambda^4}{m_c^4}\right)\sin^2\theta_C\cos^2\theta_C \times \left\{\frac{4C_1}{3}[3C_1+2C_2]B_D - \left[C_2^2 + \frac{8C_1C_2}{3} + 4C_1^2\right]\bar{B}_D^{(S)}\right\} \quad (4)$$

where Λ is a non-perturbative parameter of the order of hadronization scale or smaller (or $\Lambda \sim 1$ GeV or smaller) that is related to the light quark-antiquark condensates. It is defined as

$$m_s^2\Lambda^4 = \langle 0 | \left[(\bar{s}(0)\gamma_\mu P_L s(0) - \bar{d}(0)\gamma_\mu P_L d(0)) (\bar{s}(0)\gamma^\mu P_L s(0) - \bar{d}(0)\gamma^\mu P_L d(0)) - 2\bar{s}(0)\gamma_\mu P_L d(0)\bar{d}(0)\gamma^\mu P_L s(0) \right] | 0 \rangle \quad (5)$$

Choosing $\Lambda = 1$ GeV, one gets

$$(x_D)^{d=12}_{external} = -0.76 \times 10^{-3} \quad (6)$$

While the negative sign may be flipped when evaluating the remaining diagrams with hanging quarks, it is apparent that our result is somewhat smaller than the estimate discussed above. This is mainly due to the fact that the loop factor $1/(3\pi^2)$ present in Eq. (3) does not just disappear (as was assumed in ref. [8]) but is replaced by another suppression factor, $2/9$, due to Fierz re-arrangements of the relevant color indices when deriving (4). More detailed analysis of the power $1/N_{color}$ suppression effects is going to be presented in [15].

Thus, our quantitative evaluation of the non-perturbative contribution to the mass difference in $D^0 - \bar{D}^0$ mixing yields the value of x_D that is few times smaller than

the made estimates and by order of magnitude less than the experimental output for this quantity. The interpretation of our result is rather ambiguous. As mentioned above, it may be due to the failure of Operator Product Expansion (as μ_{had}/m_c is not a small quantity). Another issue may be the validity of the factorization approach used to derive (4). On the other hand, it is also possible that the SM contribution to x_D is indeed small (provided strong cancelations in the contribution of the exclusive channels occur), and the experimental value of x_D is due to New Physics effects.

In conclusion we attempted to evaluate quantitatively the non-perturbative contribution to the mass difference in $D^0 - \bar{D}^0$ mixing within the Standard Model using the OPE techniques. Our preliminary results show that within this approach the SM prediction for x_D is few times less than the previously made estimate and order of magnitude less than the experimental value of this quantity.

ACKNOWLEDGEMENTS

Author is grateful to the organizers of Charm-2012 workshop at University of Hawaii and the Center of Scholarly and Creative Excellence at Grand Valley State University for financial support, without which this presentation would not be possible.

This work has been partially supported by the grant DOE DE-FGO2-96ER41005.

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